

particle interaction. At the point B, which is separated from the equilibrium position by a distance $x_{of\ x_2}$, all the kinetic energy E_k passes into the potential energy, as a result of which the latter again increases by the value $U(x) = E_k$ and becomes equal to $-(U_0 - U(x))$.

If particle 2 performed purely harmonic vibrations, then following Hooke's law, the force $f(x)$ that occurs when it deviates from the equilibrium position by a distance x would be strictly proportional to this deviation and directed to the equilibrium position

$$f = -\beta x \tag{1}$$

The change in the potential energy $U(x)$ of the particle would be described by the parabola a'a'bc' (Fig. 1), the equation of which is

$$U(x) = \beta x^2/2 \tag{2}$$

This parabola is symmetric with respect to a straight line bd parallel to the U-axis and at a distance r_0 from it_0 . Therefore, the deviations x_1 and x_2 would be the same in magnitude, and the midpoint of the oscillation range A'B' would coincide with the equilibrium position 0. Heating the body in this case could not cause its expansion, since with an increase in temperature there would only be an increase in the amplitude of particle vibrations, and the average distances between them would remain unchanged.

In reality, the potential curve abc is, as can be seen from Figure 1, non-symmetric with respect to the line bd. This means that the vibrations of particles in a solid are anharmonic. To take into account the asymmetry of the potential curve, it is necessary (2) to introduce an additional term in equation (2) – $gx^3/3$, which expresses this asymmetry. Then equations (1) and (2) take the form:

$$U(x) = \frac{\beta x^2}{2} - \frac{gx^3}{3} \tag{3}$$

$$f(x) = -\frac{\partial U}{\partial x} = -\beta x + gx^2 \tag{4}$$

The unsymmetric nature of the potential curve leads to the fact that the deviations of particle 2 to the right and to the left are not the same ($x_{x_2} > x_{x_1}$). As a result, the average position of particle 2 (point O') no longer coincides with the equilibrium position (point O), but is shifted to the right. This corresponds to an increase in the average distance between particles by $\langle x \rangle$.

Thus, when heated, the volumetric expansion of the body is explained by an increase in the average distances between particles, and not by an increase in the amplitude of vibrations of particles near their equilibrium positions.

Let us estimate the coefficient of thermal expansion α . The average value of the force that occurs when the particle 2 is displaced from the equilibrium position is written according to (4)(3.19) by the relation

$$\langle f \rangle = -\beta \langle x \rangle + g \langle x^2 \rangle.$$

In free vibrations, $\langle f \rangle = 0$, so $I g \langle x^2 \rangle = \beta \langle x \rangle$. From here we find

$$\langle x \rangle = \frac{g \langle x^2 \rangle}{\beta} \tag{5}$$

Up to the second order of smallness, the potential energy of an oscillating particle is determined by the relation (3), and its average value is $\langle U(x) \rangle = \beta \frac{\langle x^2 \rangle}{2}$. Hence we find $\langle x^2 \rangle \approx \frac{2 \langle U(x) \rangle}{\beta}$

Substituting this formula in (3.20), we obtain for the value of the average distance between particles the expression: $\langle x \rangle = \frac{2g \langle U(x) \rangle}{\beta^2}$. In addition to the potential energy $U(x)$, the oscillating particle has a kinetic energy E_k , and $\langle U(x) \rangle = \langle E_k \rangle$. The total energy of the particle $E = \langle E_k \rangle + \langle U(x) \rangle = 2 \langle U(x) \rangle$. This allows the value $\langle x \rangle$ to be rewritten as follows: $\langle x \rangle = gE/\beta^2$.

The relative linear expansion, which is the ratio of the change in the average distance $\langle x \rangle$ between particles to the normal distance $\langle x \rangle r_0$ between them, is $\frac{\langle x \rangle}{r_0} = \frac{g}{\beta^2 r_0} E$, and the coefficient of linear expansion is defined as

$$\alpha = \frac{1}{r_0} \frac{d \langle x \rangle}{dT} = \frac{g}{\beta^2 r_0} \frac{dE}{dT} = \chi C_V \tag{5),}$$

Where $\chi = \frac{g}{\beta^2 r_0}$, a $C_V = \frac{dE}{dT}$ is the heat capacity assigned to a single particle.

Thus, the coefficient of linear expansion α is proportional to the heat capacity of the body.

1) In the region of high temperatures, the energy of linearly oscillating particles is equal to $E_k = kT$, the heat capacity assigned to one particle теплоёмкость $W_{\nu} = k$, so the expansion coefficient of a linear chain of atoms will be equal to

$$\alpha = \chi C_V = \frac{gk}{\beta^2 r_0} \quad (6)$$

Substituting the numerical values $g, k, \beta,$ and r_0 for various solids gives for α a value of the order of 10^{-4} - 10^{-5} , which is in satisfactory agreement with the experiment. Experience also confirms that in the region of high temperatures α is practically independent of temperature.

2) In the region of low temperatures, α behaves similarly to C_V : it decreases with decreasing temperature and tends to zero as it approaches absolute zero. In conclusion, we note that a formula similar to (6) was first proposed for metals by Gruneisen and had the form

$$\alpha = \gamma \cdot \eta \cdot C_V / 3V$$

where η is the compressibility coefficient of the metal, V is the atomic volume, and γ is the Gruneisen constant, which varies from 1.5 to 2.5 for different metals.

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